Computer simulation of the mechanisms and energetics of circumambulatory rearrangements of Dewar furan, thiophene and selenophene

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By *ab initio* [MP2(fc)/6-311+G**] and DFT (B3LYP/6-311+G**) calculations, energy barriers to thermal walk rearrangements of the Dewar type isomers of the aromatic five-membered heterocycles were found in the range of 16–22 kcal mol⁻¹.

Circumambulatory (merry-go-round) rearrangements due to the rapid displacement of transition metal- or main group elementcentered migrants around the periphery of a cyclopolyenyl ring, in particular, those related to cyclopentadiene derivatives, represent a comprehensively explored area of the dynamic stereochemistry of structurally non-rigid compounds.^{1,2} The principal mechanism that governs the fluxional behaviour of these compounds is controlled by [1,j]-sigmatropic shifts of a migrating group to the adjacent position of the ring. For the derivatives of electrically neutral even-membered ring cyclopolyenyl systems, this mechanism consists in the rearrangements of bicyclic structures, e.g., bicyclo[2.1.0]pent-2-enes and bicyclo[4.1.0]hepta-2,4-dienes (norcaradienes). The rearrangements of this type, which are also classified as walk rearrangements,3 are usually associated with the energy barriers exceeding those $(\leq 20-25 \text{ kcal mol}^{-1})$ measurable with the aid of the dynamic NMR technique. An exception to this tendency is the very low free energy barrier (~7 kcal mol⁻¹ at -136 °C)⁴ found for the circumambulatory rearrangement of the S-oxide of a tetra(trifluoromethyl) derivative of Dewar thiophene, whereas an analogous degenerate rearrangement of parent compound 1 (R = H)requires overcoming a free energy barrier of 22.1 kcal mol⁻¹ (at 157 °C).3

Walk rearrangements leading to the scrambling of ring atoms are considered as an important mechanistic route of the photochemical isomerizations of cyclopentadienes and five-membered heterocycles.⁵ At the same time, no additional experimental data are currently available on the mechanism and energetics of thermal walk rearrangements of the bicyclic structures derived from parent five-membered heterocycles: furan, thiophene and selenophene. The only computational studies of these rearrangements are restricted to the early PRDDO calculations⁶ of Dewar's pyrrole and furan. These calculations predicted very high energy barriers (68 and 76 kcal mol⁻¹, respectively) to the displacements of heteroatoms along the four-membered ring. Here we report on the *ab initio* [MP2(fc)/6-311+G**] and DFT (B3LYP/6-311+G**) theoretical study of the critical parts of the ground state potential energy surfaces (PESs) related to the degenerate thermally induced circumambulatory (walk) rearrangements of Dewar furan (X = O), thiophene (X = S) and selenophene (X = Se) 2.

Main results of the calculations are given in Table 1 and Figures 1–4. The geometries of the ground state structures of bicyclic compounds ${\bf 2}$ are shown in Figure 1. The values of $E_{\rm rel}$ in Table 1 indicate that Dewar isomers ${\bf 2}$ of the aromatic five-

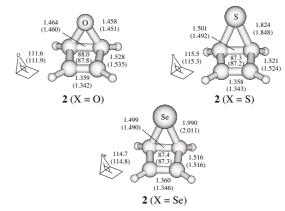


Figure 1 MP2(fc)/6-311+ G^{**} and B3LYP/6-311+ G^{**} (in parentheses) optimised geometries (bond lengths in Angström units and angles in degrees) of Dewar furan **2** (X = O), thiophene **2** (X = S) and selenophene **2** (X = Se).

membered heterocycles are considerably strained compared to the latter, with the strain progressively decreased with the size of the heteroatom. The calculations predict that all compounds **2** are susceptible to sufficiently low-energy barrier walk rearrangements occurring as the sequence of [1,3]-sigmatropic shifts around the periphery of the four-membered cyclobutene rings. The geometries of the transition state structures corresponding to a step in these rearrangements are shown in Figure 2.

The topology of the PESs of the walk rearrangements of 2, which is generally common to all the compounds, is portrayed by Figure 3. The minimal energy reaction paths for the rearrangements are stretched along the slopes of the rather high

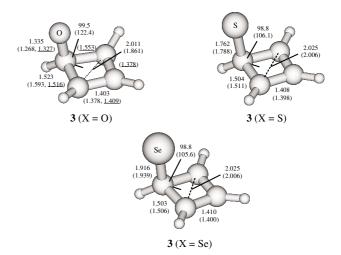


Figure 2 MP2(fc)/6-311+G** and B3LYP/6-311+G** (in parentheses) optimised geometries (bond lengths in Angström units and angles in degrees) of the transition state structures for the thermal walk rearrangements of Dewar furan **3** (X = O), thiophene **3** (X = S) and selenophene **3** (X = Se). Structural parameters of the intermediate corresponding to a shallow minimum on the B3LYP/6-311+G** PES in the vicinity of the transition state structure are given underlined.

Table 1 Total energies plus zero point energy (ZPE) corrections, relative energies of furan, thiophene, selenophene, their Dewar type isomers **2**, transition state structures **3** and stationary points **4** corresponding to a top on the hill at the PES (Figure 3) for walk rearrangements of **2** calculated by MP2(fc)/6-311+G** and B3LYP/6-311+G** methods.

Structure	$-(E_{\text{tot}} + \text{ZPE}^a)$ (a.u.)		$E_{\rm rel}/{\rm kcal~mol^{-1}}$		$\lambda^b (\mathrm{i} \omega/\mathrm{cm}^{-1})$	
	MP2(fc)/6-311+G**	B3LYP/6-311+G**	MP2(fc)/6-311+G**	B3LYP/6-311+G**	MP2(fc)/6-311+G**	B3LYP/6-311+G**
Furan	229.36820	230.01950	-65.6	-67.8	0	0
2(X = O)	229.26371	229.91149	0	0	0	0
3(X = O)	229.22947	229.89216	20.5	22.1	0^c	1 (i210)
4(X = O)	229.03974	229.77247	140.5	87.2	2	3
Thiophene	551.97962	553.00808	-57.3	-58.5	0	0
2(X = S)	551.88833	552.91483	0	0	0	0
3(X = S)	551.84635	552.88990	16.3	15.7	1 (i594)	1 (i294)
4(X = S)	551.76674	552.79888	76.3	72.8	2	2
Selenophene	2554.22535	2556.33600	-40.6	-52.5	0	0
2(X = Se)	2554.16068	2556.33601	0	0	0	0
3 (X = Se)	2554.11052	2556.23201	31.5	22.8	1 (i556)	1 (i288)
4(X = Se)	2554.03849	2556.15234	76.7	62.8	2	2

^aZero point vibration energy factors for MP2 (0.9748) and DFT (0.9806) calculations are taken according to ref. 7. ^bThe number of negative eigenvalues of the Hessian matrix. ^cAn intermediate corresponding to a very shallow minimum on the PES.

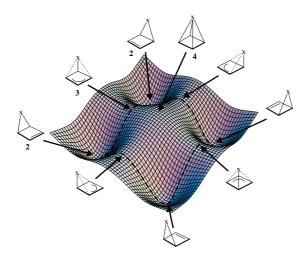


Figure 3 The pattern of the lowest singlet PESs of the walk rearrangements of Dewar furan (X = O), thiophene (X = S) and selenophene (X = Se). The minimal energy reaction path is indicated by a dashed line.

(see Table 1) hill with square-pyramidal $C_{\rm 4v}$ structures 4 on its top. Structures 4 may be considered as $D_{\rm 4h}$ cyclobutadiene capped by an oxygen (X = O) or a chalcogen (X = S, Se) atom. It was found that stable structures of the $C_{\rm nv}$ type must contain

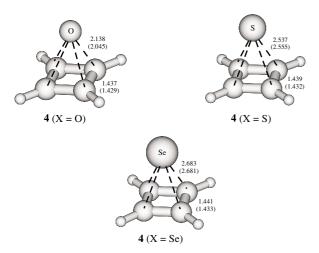


Figure 4 MP2(fc)/6-311+ G^{**} and B3LYP/6-311+ G^{**} (in parentheses) of the optimised geometries (bond lengths in Angström units) of unstable structures **4** corresponding to a top on the hills on the PESs for the walk rearrangements of Dewar furan (X = O), thiophene (X = S) and selenophene (X = Se) shown in Figure 3.

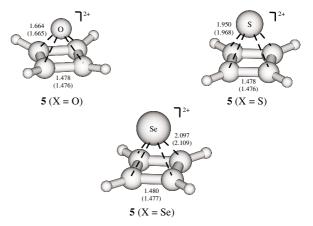


Figure 5 MP2(fc)/6-311+G** and B3LYP/6-311+G** (in parentheses) optimised geometries (bond lengths in Angström units) of square-pyramidal dications **5**.

no more than eight electrons (8e rule) in the valence orbitals of their fragments: π -MOs of the basal ring and valence AOs of the apical atom or σ -bonded group. Thus, the instability of structures **4** is caused by an excessive number of electrons (10) on these orbitals. In accordance with the 8e rule, the removal of two electrons from the MOs of **4** may lead to the formation of stable $C_4H_4X^{2+}$ dications **5**. This expectation was previously tested by semiempirical calculations 8,9 and verified by *ab initio* and DFT calculations carried out in this work. All the C_{4v} dications correspond to sufficiently deep local minima on the respective PESs. The calculated geometries of structures **4** and **5** are shown in Figures 4 and 5, respectively. The essential difference in the geometries of unstable C_{4v} structures **4** and stable C_{4v} dications **5** consists in the substantial (0.4–0.6 Å) shortening of the C–X distances in the latter, which are ~0.2 Å longer than the standard C_{-v2} –X bonds.

longer than the standard C_{sp2} –X bonds. All the calculations were performed using the Gaussian 98¹⁰ program with the 6-311+G** basis set. Electron correlation was accounted for in the second order Møller–Plesett perturbation theory approximation. In the DFT calculations, the B3LYP exchange-correlation functional was employed. All the stationary points in the PES were identified by calculations of the force constant matrix (Hessian). Zero point energies (ZPE) were corrected⁷ by multiplying them by 0.9748 [MP2(fc)/6-311+G** level] and 0.9806 (B3LYP/6-311+G** level).

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